

# Ultra-Wide-Bandgap Aluminum Gallium Nitride Devices for Radio-Frequency Applications

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**Abstract**—The motivation for the use of ultra-wide-bandgap, Al-rich AlGa<sub>N</sub> for radio-frequency high electron mobility transistors is presented. An 8× increase in RF power is expected based on the anticipated 4× increase in breakdown electric field, 2× increase in channel sheet carrier density, and parity in electron saturation velocity for Al-rich AlGa<sub>N</sub> compared to Ga<sub>N</sub>. Numerical simulations of HEMTs based on Al-rich AlGa<sub>N</sub>/AlGa<sub>N</sub> heterostructures indicate that a power density approaching 20 W/mm may be possible. Initial experimental results on an Al<sub>0.85</sub>Ga<sub>0.15</sub>N/Al<sub>0.70</sub>Ga<sub>0.30</sub>N HEMT demonstrate the viability of the technology.

**Keywords**—Ultra-Wide-Bandgap, Aluminum Gallium Nitride, High Electron Mobility Transistor, Radio Frequency

## I. INTRODUCTION

Today's state-of-the-art technology in radio-frequency (RF) power amplifiers is capable of 5 W/mm using Al<sub>x</sub>Ga<sub>1-x</sub>N/GaN HEMTs, with  $x$  typically in the range 0.25 – 0.30. The emerging material system based on Ultra-Wide-Bandgap (UWBG) Al-rich Al<sub>x</sub>Ga<sub>1-x</sub>N/Al<sub>y</sub>Ga<sub>1-y</sub>N heterostructures with  $0.7 \leq y < x \leq 1.0$  has the potential to greatly exceed the capabilities of today's AlGa<sub>N</sub>/Ga<sub>N</sub> HEMTs, possibly achieving 20 W/mm power density. This projection stems from the expected 4× enhancement of critical electric field, 2× enhancement of sheet carrier density, and parity of electron saturation velocity for Al-rich AlGa<sub>N</sub> channel HEMTs relative to Ga<sub>N</sub>-channel HEMTs. The anticipated increase in RF power density for Al-rich AlGa<sub>N</sub>-channel HEMTs is due in part to similarities to the existing AlGa<sub>N</sub>/Ga<sub>N</sub> system, in part to theoretical analysis, and in part to existing data for long-channel Al-rich AlGa<sub>N</sub>-channel devices for power switching applications [1-5]. This work presents the theoretical basis for the use of Al-rich AlGa<sub>N</sub> for RF applications, and also presents initial experimental results for Al<sub>0.85</sub>Ga<sub>0.15</sub>N/Al<sub>0.70</sub>Ga<sub>0.30</sub>N RF HEMTs.

## II. MOTIVATION FOR AL-RICH ALGAN RF DEVICES

To first order, the maximum RF power output from a transistor is given by the equation:

$$P_{max} = I_{max} V_{max} / 8 \quad (1)$$

where  $I_{max}$  is the maximum RF current and  $V_{max}$  is the difference between the HEMT breakdown voltage and the knee voltage. Al-rich AlGa<sub>N</sub>-channel HEMTs have the potential in identically-sized devices to increase  $P_{max}$  relative to Ga<sub>N</sub>-channel HEMTs due to a larger critical electric field  $E_c$  resulting from a larger bandgap  $E_g$ . The critical electric field of AlGa<sub>N</sub> has been observed to obey the power-law scaling:

$$E_c \sim E_g^n \quad (2)$$

with  $n$  in the range 2.5 – 2.7 [6], consistent with the prediction of Hudgins [7]. Current work is re-examining the reported experimental critical field values and their dependence on bandgap in more detail, and is also attempting to provide a theoretical explanation for the empirically observed power-law relationship between critical field and bandgap. A practical Ga<sub>N</sub>-channel HEMT operates at an effective breakdown electric field of approximately 100 V/μm, roughly a quarter of the theoretical value. Assuming a similar 4× theoretical-to-achievable reduction factor for Al-rich AlGa<sub>N</sub>-channel HEMTs, it is expected that such a device could achieve a breakdown field of approximately 400 V/μm, or roughly 4× greater than that of a Ga<sub>N</sub>-channel HEMT, for alloy compositions  $x$  in the 0.70 to 1.00 range. The significance of the 4× higher critical electric field for an Al-rich AlGa<sub>N</sub>-channel HEMT is that the sub-micron channel dimensions required for RF operation allow for biasing at a much larger voltage than is possible for a Ga<sub>N</sub>-channel HEMT of the same physical size. Thus, a higher breakdown voltage for equivalent device dimensions (i.e. equivalent cutoff frequency  $f_T$ ) should be achievable for AlGa<sub>N</sub>-channel HEMTs relative to Ga<sub>N</sub>-channel HEMTs. Since power output in an RF power amplifier is proportional to the product of  $I_{max}$  and  $V_{max}$  (Equation 1) the 4× higher critical electric field for AlGa<sub>N</sub>-channel HEMTs is expected to enable increases in power output deriving from higher voltage operation, as long as the requisite high current density can be achieved.

Regarding the current density, if the electric field is sufficiently high such that the velocity is saturated (as will be the case for RF operation), the current density is given by

$$J = q v_{sat} n_s \quad (3)$$

where  $q$  is the electron charge,  $v_{sat}$  is the electron saturation velocity, and  $n_s$  is the channel sheet carrier density. Based on Monte-Carlo simulations of electron saturation velocity in AlGa<sub>N</sub>, the saturation velocity is predicted to be of order  $2 \times 10^7$  cm/s over the entire composition range, so that  $v_{sat}$  for Al-rich AlGa<sub>N</sub> is expected to be roughly equivalent to that of Ga<sub>N</sub> [8]. Initial experimental measurements of Al<sub>0.70</sub>Ga<sub>0.30</sub>N by our group have yielded  $v_{sat} \approx 4 \times 10^6$  cm/s [9], and additional experiments are underway. From Equation (3), the other factor that determines current density,  $n_s$ , has already been observed by our group to exceed  $2 \times 10^{13}$  cm<sup>-2</sup>, roughly double that typically observed for traditional AlGa<sub>N</sub>/Ga<sub>N</sub> heterostructures. Indeed, there are good theoretical reasons to believe that achieving higher  $n_s$  in AlGa<sub>N</sub>-channel HEMTs than in Ga<sub>N</sub>-channel HEMTs is plausible. One reason for this is the ability to tolerate a higher bandgap offset between channel and barrier, due to the higher strength of the Al-N bond compared to the Ga-N bond. Moreover, since the polarization charge increases

with increasing Al composition, at higher Al composition higher sheet carrier density is expected for comparable composition offset. Finally, since Al-rich heterostructures would likely be grown on AlN substrates, the barrier layer would be grown under compressive strain and could thus be grown thicker than the AlGa<sub>N</sub> barrier in an AlGa<sub>N</sub>/Ga<sub>N</sub> HEMT. Due to these factors, a 2× improvement in  $n_s$  is expected for Al-rich heterostructures compared to conventional AlGa<sub>N</sub>/Ga<sub>N</sub> heterostructures. Provided that the electric field is sufficiently high such that saturation velocity is achieved, this will result in a 2× increase in current density for an Al-rich AlGa<sub>N</sub>-channel device relative to a Ga<sub>N</sub>-channel device, assuming parity in  $v_{sat}$ . All told, the expected 2× improvement in  $I_{max}$ , coupled with the 4× enhancement anticipated for  $V_{max}$ , predicts an 8× improvement in  $P_{max}$  for an Al-rich AlGa<sub>N</sub>-channel RF HEMT compared to a conventional Ga<sub>N</sub>-channel RF HEMT.

Other aspects of the AlGa<sub>N</sub> material system may also lead to improved RF performance. The knee voltage is important because  $V_{max} = V_{br} - V_{knee}$ , so that  $V_{knee}$  impacts  $P_{max}$ . Ga<sub>N</sub>-channel HEMTs have higher channel electron mobility than AlGa<sub>N</sub>-channel HEMTs do, and thus possess a smaller knee voltage. However, as a fraction of  $V_{max}$  this benefit is minimal at room temperature. Furthermore, RF devices typically operate at elevated temperatures. Since the low-field mobility of AlGa<sub>N</sub> is dominated by alloy scattering, it is not expected to be significantly less at elevated temperature than at room temperature [10], while the mobility of Ga<sub>N</sub> will degrade considerably at high temperatures since it is dominated by phonon scattering. Thus, at high temperatures the knee voltage for a Ga<sub>N</sub>-channel device is expected to be only slightly better than that of an AlGa<sub>N</sub>-channel device. Moreover, as a fraction of  $V_{max}$  the knee voltage of a Ga<sub>N</sub>-channel device will likely be relatively worse than that of an AlGa<sub>N</sub>-channel device. These considerations suggest that AlGa<sub>N</sub>-channel devices not only offer potentially better power density than Ga<sub>N</sub>-channel devices, but also superior Power-Added Efficiency (PAE). Finally, they may also offer improved linearity. Because RF linearity is lost at high power due to mobility-degradation-related channel resistance increases that result in a flattening of output power, AlGa<sub>N</sub>-channel devices are expected to show more robust linearity due to the insensitivity of mobility to temperature as discussed above.

### III. DESIGN EXAMPLE

As an example, a HEMT consisting of an AlN barrier and an Al<sub>0.70</sub>Ga<sub>0.30</sub>N channel was simulated [11]. A schematic drawing of the structure is shown in Figure 1. The barrier thickness underneath the gate, source, and drain metallizations was 10 nm, and the barrier thickness in the source-to-gate and gate-to-drain access regions was varied between 10 and 60 nm. The source-to-gate distance was 100 nm, the gate length was 100 nm, and the gate-to-drain distance was 500 nm. An Al<sub>0.78</sub>Ga<sub>0.22</sub>N back-barrier was included in the design to reduce leakage current by improving carrier confinement, and the

channel layer thickness between the barrier and the back-barrier was 150 nm.

Plots of the simulated drain current ( $I_d$ ) and transconductance ( $G_m$ ) vs. gate voltage are shown in Figure 2 for a drain-to-source voltage of 40 V and four values of the barrier thickness in the access regions. It is apparent that the thicker barrier results in not only a higher drain current, but also a higher, broader, and flatter transconductance curve, which are desirable attributes for higher power and greater linearity. However, the thicker barrier also results in higher drain leakage current in pinch-off due to the increased carrier density in the channel. For this reason, gain vs. frequency curves were simulated for the 40-nm-thick-barrier device. Mason's Unilateral Gain (MUG) vs. frequency curves for various gate voltages, derived from small-signal AC simulations, are shown in Figure 3. Biasing the device near the point of maximum transconductance ( $V_{gs} = 1$  V) results in a gain of 34 dB at 30 GHz. Further,  $f_T \approx 46$  GHz and  $f_{max} \approx 230$  GHz are predicted, based on simulated  $|h_{21}|$  vs. frequency curves (not shown).

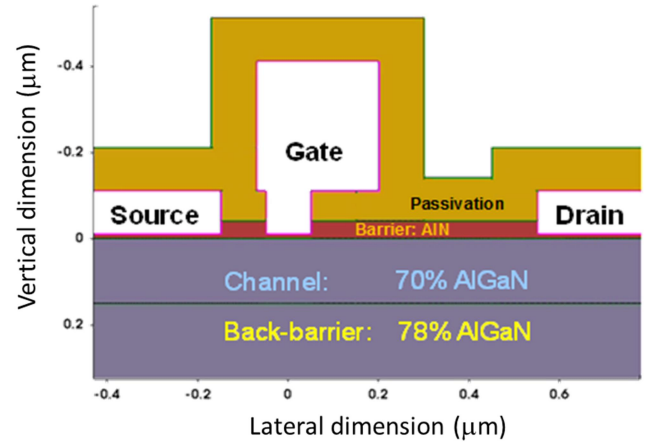


Figure 1. Schematic drawing of simulated AlN/Al<sub>0.70</sub>Ga<sub>0.30</sub>N HEMT.

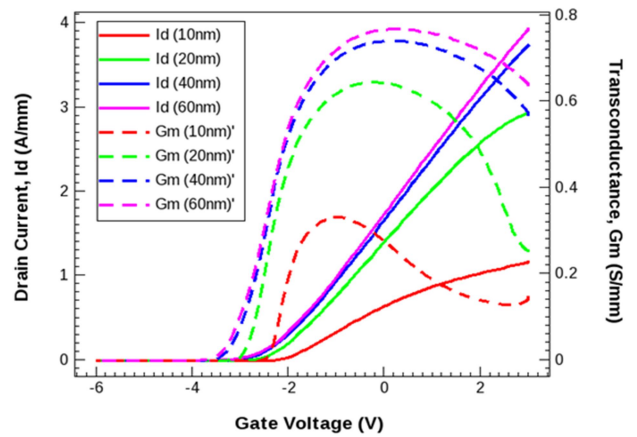


Figure 2. Simulated drain current and transconductance curves plotted vs. gate voltage for AlN/Al<sub>0.70</sub>Ga<sub>0.30</sub>N HEMT for four access region barrier thicknesses. The drain voltage is 40 V.

The simulated drain current vs. drain voltage curves are shown in Figure 4. The bias current determines the gain, PAE, and third-order intercept point IP3. The best IP3 performance will be obtained at a DC bias (Q-point) near the maximum  $G_m$

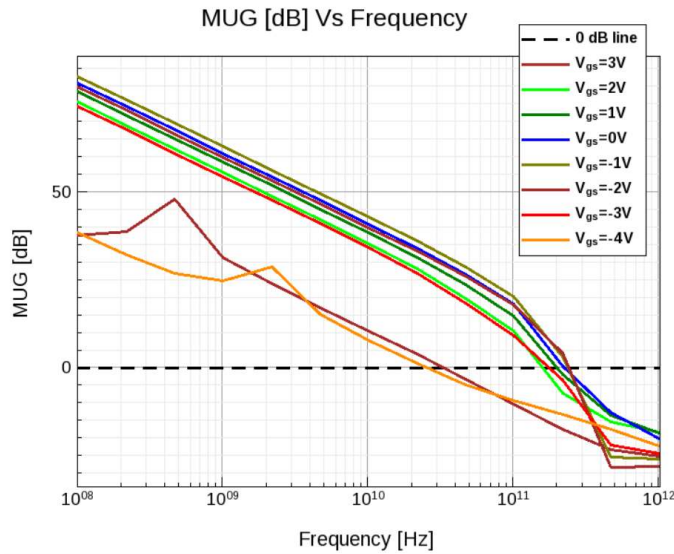


Figure 3. Simulated MUG for the 40-nm-thick barrier AlN/Al<sub>0.70</sub>Ga<sub>0.30</sub>N device. The drain voltage is 40 V.

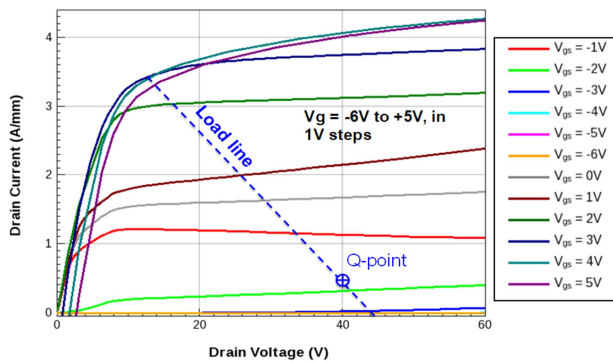


Figure 4. Drain current vs. drain voltage curves for the simulated device, with a resistive load-line superimposed.

point, which is at a drain current of  $\sim 1.8$  A/mm (Figure 2). However, PAE will suffer because of the high DC bias current. Alternatively, a lower bias current will provide higher PAE, but IP3 performance will be degraded. To obtain first-order power and linearity performance, a Q-point of  $V_{gs} = -1.75$  V,  $I_d = 0.37$  A/mm was selected, even though this Q-point is not optimum for IP3.

Large-signal DC sweeps were performed using a load-line connecting the selected Q-point to the knee voltage at  $V_{ds} \approx 10$  V,  $I_d = 3.4$  A/mm, resulting in a load resistance of  $9.1 \Omega\text{-mm}$  (Figure 4). The  $V_{out}$  vs.  $V_{in}$  transfer curve was extracted by sweeping the gate input voltage (Figure 5a). From the extracted transfer curve, a MATLAB-based algorithm was used to calculate the output voltage and power waveforms for a 30 GHz sinusoidal input (Figure 5b). The input power was calculated from  $V_{in}^2 / \text{Re}[Z_{in}]$ , where  $Z_{in}$  is the input impedance at the Q-point at 30 GHz. The output power at the fundamental (30 GHz) and the 3<sup>rd</sup> harmonic (90 GHz) was calculated from

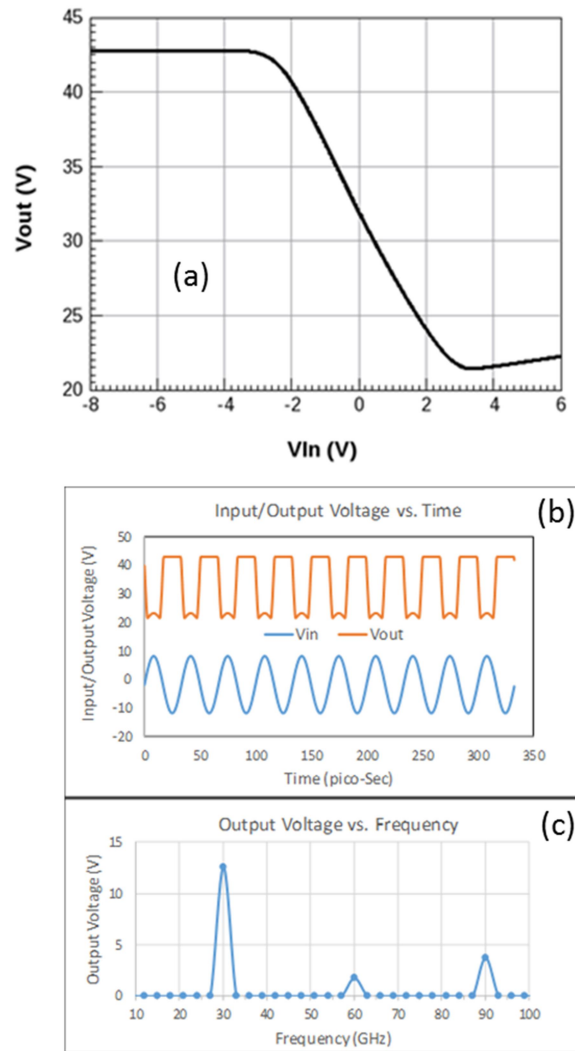


Figure 5. For simulated AlGaIn-channel HEMT: (a)  $V_{out}$  vs.  $V_{in}$  transfer curve; (b) Input and output voltage waveforms; and (c) Frequency spectrum of the output voltage at 30 dBm input power.

the Fourier transform of the output voltage waveforms (Figure 5c). PAE was calculated from the input and output powers at 30 GHz and the DC power dissipation.

The calculated power performance for the AlGaIn-channel HEMT is shown in Figure 6. For this design, a power density of 18 W/mm is achieved, substantially surpassing today's state-of-the-art  $\sim 5$  W/mm. Several approaches exist that could plausibly increase the power density, for example increasing the barrier thickness, decreasing the Al mole fraction in the channel, and lowering the Al mole fraction in the back-barrier. All these approaches will result in increased charge and current density in the channel. A thicker barrier will also reduce the peak electric field within the device, likely resulting in not only a higher maximum voltage, but also in improved reliability.

#### IV. INITIAL EXPERIMENTAL DATA

$\text{Al}_{0.85}\text{Ga}_{0.15}\text{N}/\text{Al}_{0.70}\text{Ga}_{0.30}\text{N}$  RF HEMTs were grown, fabricated, and characterized [12]. The growth was done by Metal-Organic Chemical Vapor Deposition (MOCVD) on sapphire substrates, and devices were processed in a manner similar to previous work on long-channel devices by our group [2,3]. The HEMTs were constructed using conventional rectangular geometry, a gate length of 80 nm, source-to-gate spacing of 0.5  $\mu\text{m}$ , and gate-to-drain spacing of 2.0  $\mu\text{m}$ . DC characterization of 50- $\mu\text{m}$ -wide devices (Figure 7) yielded  $\sim 80$  mA/mm at the bias condition  $V_{gs} = 5$  V,  $V_{ds} = 10$  V, with peak transconductance of 24 mS/mm near this bias point. Drain current increased to nearly 160 mA/mm when the device was biased at  $V_{gs} = 10$  V,  $V_{ds} = 10$  V.

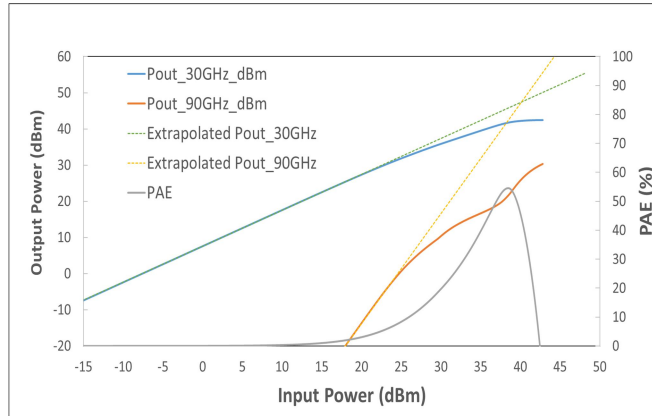


Figure 6. Power performance of the simulated AlGaIn-channel HEMT.

Current and power gain vs. frequency plots derived from S-parameter measurements were used to determine  $f_T$  and  $f_{max}$  as a function of drain voltage. Results are plotted in Figure 8 as a function of drain voltage, at a gate voltage  $V_{gs} = 2.75$  V. At the bias condition  $V_{gs} = 2.75$  V,  $V_{ds} = 20$  V,  $f_T \approx 28.4$  GHz and  $f_{max} \approx 18.5$  GHz. The reversal in order of frequency from what is expected (i.e.  $f_T > f_{max}$  is observed, while typically  $f_{max} > f_T$ ) may be due to gate resistance and/or drain conductance. Large-signal operation of a two-finger device at 3 GHz was performed at  $V_{gs} = 3.75$  V,  $V_{ds} = 20$  V to approximate class-A amplifier operation. Results for output power  $P_{out}$ , PAE, power

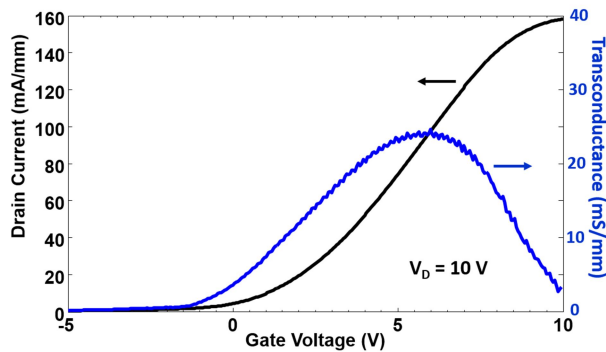


Figure 7. DC transfer characteristics for single-finger  $\text{Al}_{0.85}\text{Ga}_{0.15}\text{N}/\text{Al}_{0.70}\text{Ga}_{0.30}\text{N}$  HEMT with 80 nm gate length and 50  $\mu\text{m}$  gate width.

gain  $G_p$ , and transducer gain  $G_t$  are plotted as a function of input power  $P_{in}$  in Figure 9.  $P_{out} = 15.8$  dBm (0.38 W/mm) at the maximum PAE of 11%. These experimental results fall short of the simulated results, but are expected to improve considerably with improvements to the Ohmic contacts. These initial results demonstrate the viability of Al-rich AlGaIn HEMT technology for RF applications, and compare favorably to other emerging UWBG technologies in this arena, such as  $\text{Ga}_2\text{O}_3$  [13].

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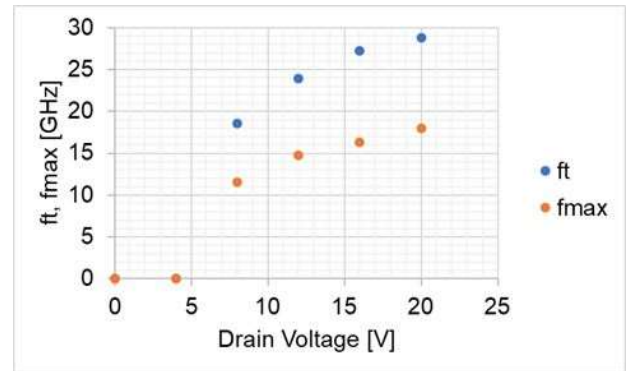


Figure 8. Cutoff frequency  $f_T$  and maximum oscillating frequency  $f_{max}$  for the HEMT of Figure 7 measured vs. drain voltage at  $V_{gs} = 2.75$  V.

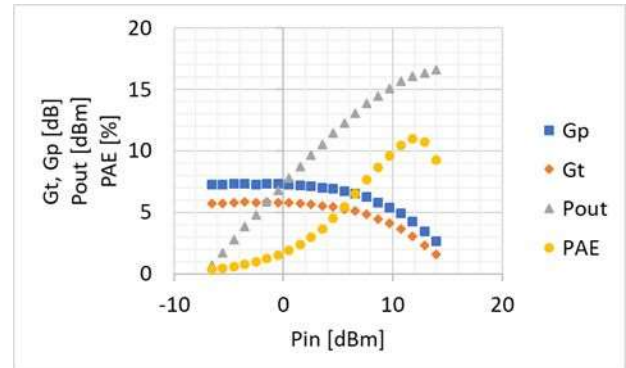


Figure 9.  $P_{out}$ , PAE,  $G_p$ , and  $G_t$  vs.  $P_{in}$  for two-finger  $\text{Al}_{0.85}\text{Ga}_{0.15}\text{N}/\text{Al}_{0.70}\text{Ga}_{0.30}\text{N}$  HEMT with 80-nm-long gate.

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